

Role of non-coplanarity in nuclear reactions using the Wong formula based on the proximity potential

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Introduction

Recently [1], we assessed Wong's formula [2] for its angular momentum ℓ -summation and "barrier modification" effects at sub-barrier energies in the dominant fusion-evaporation and capture (equivalently, quasi-fission) reaction cross-sections. For use of the multipole deformations (up to β_4) and (in-plane, $\Phi=0^0$) orientations-dependent proximity potential in fusion-evaporation cross-sections of $^{58}\text{Ni}+^{58}\text{Ni}$, $^{64}\text{Ni}+^{64}\text{Ni}$ and ^{100}Mo , known for fusion hindrance phenomenon in coupled-channels calculations, and the capture cross-sections of $^{48}\text{Ca}+^{238}\text{U}$, ^{244}Pu and ^{248}Cm reactions, forming superheavy nuclei, though the simple $\ell=0$ barrier-based Wong formula is found inadequate, its extended version, the ℓ -summed Wong expression fits very well the above noted capture cross-sections at all center-of-mass energies $E_{c.m.}$'s, but require (additional) modifications of the barriers to fit the fusion-evaporation cross-sections in the Ni-based reactions at below-barrier energies. Some barrier modification effects are shown [1] to be already present in Wong expression due to its inbuilt ℓ -dependence via ℓ -summation.

In this paper, we study for the first time the dynamics of fission reactions, such as $^{11}\text{B}+^{235}\text{U}$ and $^{14}\text{N}+^{232}\text{Th}$ forming $^{246}\text{Bk}^*$ [3], on the basis of the extended, ℓ -summed Wong formula, including also the non-coplanarity ($\Phi \neq 0^0$) degree-of-freedom for all the three types of reactions, the fusion-evaporation, capture and fission cross-sections.

The extended Wong model

Wong's expression for fusion cross-section due to colliding two deformed and oriented nuclei (orientations θ_i), lying in two different planes (azimuthal angle Φ between the

planes), in terms of ℓ partial waves, is

$$\sigma(E_{c.m.}, \theta_i, \Phi) = \frac{\pi}{k^2} \sum_{\ell=0}^{\ell_{max}} (2\ell+1) P_{\ell}(E_{c.m.}, \theta_i, \Phi), \quad (1)$$

with $k = \sqrt{\frac{2\mu E_{c.m.}}{\hbar^2}}$, and μ , the reduced mass. P_{ℓ} is the transmission coefficient for each ℓ , describing, in Hill-Wheeler approximation, the penetration of barrier $V_{\ell}(R, E_{c.m.}, \theta_i, \Phi)$.

Instead of solving Eq. (1) explicitly, which require the complete ℓ -dependent potentials $V_{\ell}(R, E_{c.m.}, \theta_i, \Phi)$, Wong summed it up *approximately*, using only $\ell=0$ quantities, which on replacing the ℓ -summation in (1) by an integral, gives the Wong formula [2]

$$\begin{aligned} \sigma(E_{c.m.}, \theta_i, \Phi) &= \frac{R_B^0{}^2 \hbar \omega_0}{2E_{c.m.}} \ln \left[1 + \exp \left(\frac{2\pi}{\hbar \omega_0} (E_{c.m.} - V_B^0) \right) \right]. \end{aligned} \quad (2)$$

Integrating (2) over θ_i and Φ , we get the fusion cross-section $\sigma(E_{c.m.})$.

For an explicit summation over ℓ in Eq. (1), the ℓ -dependent interaction potential $V_{\ell}(R)$ is a sum of Coloumb and nuclear proximity and centrifugal potentials, as

$$\begin{aligned} V_{\ell}(R) &= V_P(R, A_i, \beta_{\lambda i}, T, \theta_i, \Phi) + \frac{\hbar^2 \ell(\ell+1)}{2\mu R^2} \\ &+ V_C(R, Z_i, \beta_{\lambda i}, T, \theta_i, \Phi), \end{aligned} \quad (3)$$

where, the ℓ -summation in Eq. (1) is then carried out for the ℓ_{max} determined empirically for a best fit to measured cross-section. This procedure of explicit ℓ -summation works very well for $\Phi=0^0$ case [1] in capture reactions $^{48}\text{Ca}+^{238}\text{U}$, ^{244}Pu and ^{248}Cm , but require further modification of the barrier for Ni-based reactions at sub-barrier energies, which could be carry out empirically [1] by either (i) keeping the curvature $\hbar\omega_{\ell}$ same and modifying the

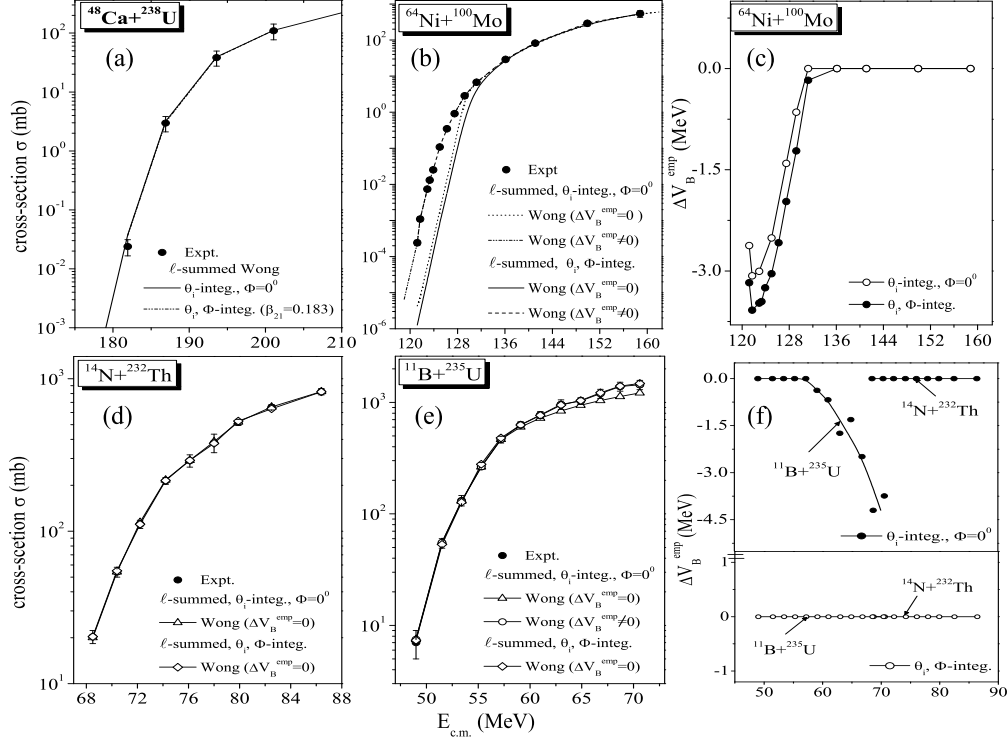


FIG. 1: ℓ -summed Wong results for both $\Phi=0^0$ and $\Phi \neq 0^0$, compared with experimental data, for: (a) $^{48}\text{Ca}+^{238}\text{U}$ (b) and (c) $^{64}\text{Ni}+^{100}\text{Mo}$, and (d) to (f) $^{246}\text{Bk}^*$ due to $^{14}\text{N}+^{232}\text{Th}$ and $^{11}\text{B}+^{235}\text{U}$ channels.

barrier height V_B^ℓ , as

$$V_B^\ell(\text{modified}) = V_B^\ell + \Delta V_B^{\text{emp}},$$

or (ii) keep the barrier height V_B^ℓ same and modify the curvature $\hbar\omega_\ell$. We use here the method of modifying the barrier height.

Calculations and results

The results of ℓ -summed Wong expression (1) for both the cases of $\Phi=0^0$ and $\Phi \neq 0^0$ are given in Fig. 1 for all the three types of reactions. Fig. 1(a) shows that the capture cross-section in $^{48}\text{Ca}+^{238}\text{U}$ is fitted nicely even after giving a small deformation ($\beta_{21}=0.183$) to ^{48}Ca for carrying out Φ -integration. The fitted $\ell_{\text{max}}(E_{\text{c.m.}})$ increase by one-to-two units. Similarly, the $^{64}\text{Ni}+^{100}\text{Mo}$ reaction is fitted for both $\Phi=0^0$ and $\Phi \neq 0^0$ by allowing empirically a small increase in “barrier lowering” ΔV_B^{emp} (Figs. 1(b) and 1(c)). On the other hand,

there is a strong entrance channel dependence in the case of fission reaction: whereas a nice fit is obtained for both $\Phi=0^0$ and $\Phi \neq 0^0$ cases in $^{14}\text{N}+^{232}\text{Th}$ channel (Fig. 1(d)), a large disagreement in cross-sections at higher energies (Fig. 1(e)) and hence a large “barrier lowering” ΔV_B^{emp} (Fig. 1(f)) is obtained for $\Phi=0^0$ in $^{11}\text{B}+^{235}\text{U}$ channel, which reduces to zero for $\Phi \neq 0^0$ case. In other words, for fission of $^{246}\text{Bk}^*$, the inclusion of non-coplanarity gives a complete fit to data for both the reaction channels, without introducing ΔV_B^{emp} .

References

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